

## (b) Energy Distribution of Ideal Gas in Three Dimensions

Entropy ( $S$ ) is measurement of randomness of system. It is function of number of microstate (which is number of ways to achieve any energy  $E$ ). Statistically it is seen so  $S = k_B \ln W$ . And Law of nature reveal that at equilibrium the entropy is maximum so  $dS = 0 \Rightarrow \ln W = 0$

For distinguishable particle. Taking in for equation  $W = N! \prod_{i=1}^l \left[ \frac{g_i^{n_i}}{n_i!} \right]$

$$\ln W = \ln \left[ N + \sum_{i=1}^l \left[ n_i \ln g_i - \ln n_i \right] \right]$$

Using sterling approximation  $\ln n = n \ln n - n$

$$= N \ln N - N + \sum_{i=1}^l n_i \ln g_i - \sum_{i=1}^l n_i \ln n_i + \sum_{i=1}^l n_i$$

$$d \ln W = \sum_{i=1}^l dn_i \ln g_i - \sum_{i=1}^l \left[ n_i \frac{1}{n_i} dn_i + \ln n_i dn_i \right] + \sum_{i=1}^l dn_i$$

$$dW = \sum_{i=1}^l dn_i \ln g_i - \sum_{i=1}^l \ln n_i dn_i = \sum_{i=1}^l \ln \frac{g_i}{n_i} dn_i = 0$$

$$\sum n_i = N \Rightarrow \sum dn_i = 0 \Rightarrow \alpha \sum dn_i = 0$$

$$\sum E_i n_i = U \Rightarrow \sum E_i dn_i = 0 \Rightarrow \beta \sum E_i dn_i = 0$$

where  $\alpha$  and  $\beta$  are used as LaGrange's multiplier which is using to make equation dimensionless where dimension of  $\beta$  is inverse of energy can be related to equilibrium

temperature  $T$  so  $\beta = \frac{1}{k_B T}$ .  $\alpha$  can be identified later can be related to fugacity

$$\sum_{i=1}^l \ln \frac{g_i}{n_i} dn_i = \alpha \sum_{i=1}^l dn_i + \beta \sum_{i=1}^l E_i dn_i$$

Equating the coefficient of  $dn_i$  then  $\frac{n_i}{g_i} = \exp(-\alpha) \exp(-\beta E)$

So, Maxwell distribution is given by  $f(\varepsilon) = \frac{n_i}{g_i} = e^{-\alpha} e^{-\beta E}$

## Calculation of $\alpha$ in three dimensional case

Total number of particles in ideal gas  $N = \sum_{i=1}^l n_i \Rightarrow N = \sum g_i e^{-\alpha} e^{-\beta \epsilon}$

where energy levels are continuous then  $N = e^{-\alpha} \int_0^{\infty} g(E) e^{-\beta E} dE$

for three dimensional case  $g(E) = \frac{2\pi V}{h^3} (2m)^{3/2} E^{1/2}$

$$N = e^{-\alpha} \int_0^{\infty} \frac{2\pi V}{h^3} (2m)^{3/2} E^{1/2} e^{-\beta E} dE$$

$$N = e^{-\alpha} \frac{2\pi V}{h^3} (2m)^{3/2} \int_0^{\infty} E^{1/2} e^{-\beta E} dE = e^{-\alpha} \frac{2\pi V}{h^3} (2m)^{3/2} \left(\frac{1}{\beta}\right)^{3/2} \left[\frac{3}{2}\right]$$

where  $\beta = \frac{1}{k_B T}$ ,  $\left[\frac{3}{2}\right] = \frac{1}{2} \sqrt{\pi}$

## Derivation of Maxwell-Boltzmann Distribution in Three Dimension

The Maxwell-Boltzmann distribution law for the particles in the states is

$$n_i = g_i \exp(-\alpha - \beta E_i) \quad , \quad n_i = g_i (\exp -\alpha) (\exp - \beta E_i)$$

After using the values  $e^{-\alpha} = \frac{N}{V} \left(\frac{h^2}{2\pi m k_B T}\right)^{3/2}$ , where  $\beta = \frac{1}{k_B T}$

We get  $\frac{n_i}{g_i} = f(E) = \frac{N}{V} \left(\frac{h^2}{2\pi m k_B T}\right)^{3/2} \exp\left(-\frac{E}{k_B T}\right)$

The number of particles  $dN(E)$  having energies in the range from  $E$  to  $E + dE$  is

$dN(E) = f(E)g(E)dE$  where  $f(E)$  is distribution function and  $g(E)dE$  is number of level (quantum state) in the range of  $E$  to  $E + dE$

The number of particles  $dn(E)$  having energies in the range from  $E$  to  $E + dE$  in three dimensional space

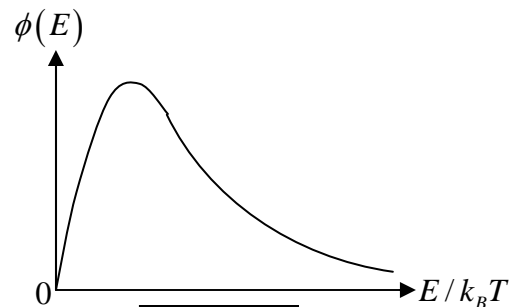


Figure 1

$$dN = \frac{N}{V} \left( \frac{h^2}{2\pi m k_B T} \right)^{3/2} \exp\left(-\frac{E}{k_B T}\right) 2\pi V \left( \frac{2m}{h^2} \right)^{3/2} E^{1/2} dE$$

$$dN = \frac{2\pi N}{(\pi k_B T)^{3/2}} E^{1/2} \exp\left(-\frac{E}{k_B T}\right) dE$$

$\frac{dN}{N}$  is fraction of particles that have energy between from  $E$  to  $E + dE$  is

$$\frac{dN}{N} = \frac{2\pi}{(\pi k_B T)^{3/2}} E^{1/2} e^{-E/k_B T} dE \text{ which is popularly known as probability to find gas have}$$

energy between from  $E$  to  $E + dE$   $\frac{dN}{N} = \phi(E) dE = \frac{2\pi}{(\pi k_B T)^{3/2}} E^{1/2} \exp\left(-\frac{E}{k_B T}\right) dE$ ,

where  $\phi(E) = \frac{2\pi}{(\pi k_B T)^{3/2}} E^{1/2} \exp\left(-\frac{E}{k_B T}\right)$  is identified probability density

This is known as the Maxwell-Boltzmann energy distribution law for an ideal gas,

where  $\lambda = \frac{h}{\sqrt{2\pi m k_B T}}$  is defined as the thermal wavelength.