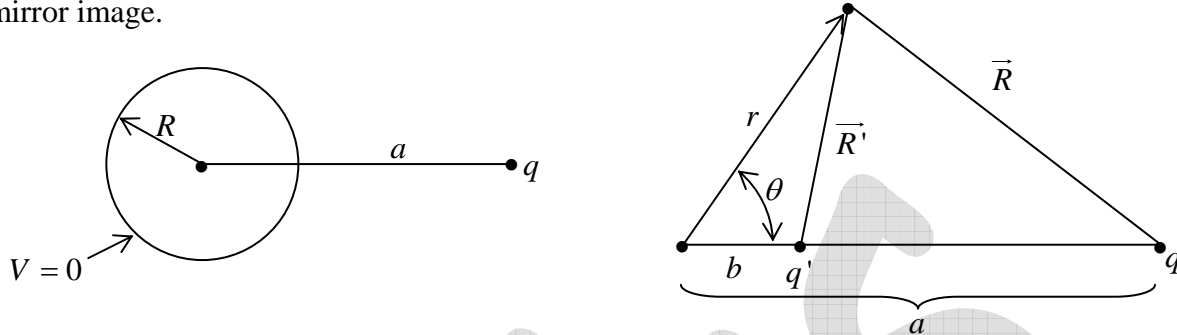


(n) Other Image Problem

The method just described is not limited to a single point charge; any stationary charge distribution near a grounded conducting plane can be treated in the same way, by introducing its mirror image.



Let us examine the completely different configuration, consisting of the point charge q together with another point charge $q' = -\frac{R}{a}q$.

Placed at a distance $b = \frac{R^2}{a}$

to the right of the centre of sphere. No conductor, now-just two point charges. The potential of

this configuration is $V(r, \theta) = \frac{1}{4\pi\epsilon_0} \left(\frac{q}{R} + \frac{q'}{R'} \right)$

where $R = \sqrt{r^2 + a^2 - 2ra \cos \theta}$, $R' = \sqrt{r^2 + b^2 - 2rb \cos \theta}$

$$\Rightarrow V(r, \theta) = \frac{1}{4\pi\epsilon_0} \left\{ \frac{q}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} + \frac{q'}{\sqrt{r^2 + b^2 - 2rb \cos \theta}} \right\}$$

$$V(r, \theta) = \frac{q}{4\pi\epsilon_0} \left\{ \frac{1}{\sqrt{r^2 + a^2 - 2ra \cos \theta}} - \frac{1}{\sqrt{R^2 + (ra/R)^2 - 2ra \cos \theta}} \right\}$$

Clearly when $r = R$, $V \rightarrow 0$

Induced charge

$\sigma = -\epsilon_0 \frac{\partial V}{\partial n}$. In this case $\frac{\partial V}{\partial n} = \frac{\partial V}{\partial r}$ at the point $r = R$.

$$\Rightarrow \sigma(\theta) = -\epsilon_0 \left. \frac{\partial V(r, \theta)}{\partial r} \right|_{r=R}$$

$$\sigma(\theta) = -\frac{q}{4\pi} \left\{ -(R^2 + a^2 - 2Ra \cos \theta)^{-3/2} (R - a \cos \theta) + (R^2 + a^2 - 2Ra \cos \theta)^{-3/2} \left(\frac{a^2}{R} - a \cos \theta \right) \right\}$$

$$\sigma(\theta) = \frac{q}{4\pi} (R^2 + a^2 - 2Ra \cos \theta)^{-3/2} \left[R - a \cos \theta - \frac{a^2}{R} + a \cos \theta \right]$$

$$\sigma(\theta) = \frac{q}{4\pi R} (R^2 - a^2) (R^2 + a^2 - 2Ra \cos \theta)^{-3/2}$$

$$q_{included} = q' = \int \sigma da = \frac{q}{4\pi R} (R^2 - a^2) \int (R^2 + a^2 - 2Ra \cos \theta)^{-3/2} R^2 \sin \theta d\theta d\phi$$

$$q' = \frac{q}{4\pi R} (R^2 - a^2) 2\pi R^2 \left[-\frac{1}{Ra} (R^2 + a^2 - 2Ra \cos \theta)^{-1/2} \right]_0^\pi$$

$$q' = \frac{q}{2a} (a^2 - R^2) \left[\frac{1}{\sqrt{R^2 + a^2 + 2Ra}} - \frac{1}{\sqrt{R^2 + a^2 - 2Ra}} \right]$$

But $a > R$ (else q would be inside), so $\sqrt{R^2 + a^2 - 2Ra} = a - R$

$$q' = \frac{q}{2a} (a^2 - R^2) \left[\frac{1}{(a+R)} - \frac{1}{(a-R)} \right] = -\frac{q}{a} R \Rightarrow q' = -\frac{q}{a} R$$

Force

The force on q , due to the sphere, is the same as the force of the image charge q' , thus:

$$F = \frac{1}{4\pi\epsilon_0} \frac{qq'}{(a-b)^2} = \frac{1}{4\pi\epsilon_0} \left(-\frac{R}{a} q^2 \right) \frac{1}{\left(a - \frac{R^2}{a} \right)^2} = -\frac{1}{4\pi\epsilon_0} \frac{q^2 Ra}{(a^2 - R^2)^2}$$

Energy

To bring q in from infinity to a , we do work

$$W = \frac{q^2 R}{4\pi\epsilon_0} \int_\infty^a \frac{\bar{a}}{(\bar{a}^2 - R^2)^2} d\bar{a} = \frac{q^2 R}{4\pi\epsilon_0} \left[-\frac{1}{2} \frac{1}{(\bar{a}^2 - R^2)} \right]_\infty^a = -\frac{1}{4\pi\epsilon_0} \frac{q^2 R}{2(a^2 - R^2)}$$

$$W = -\frac{1}{4\pi\epsilon_0} \frac{q^2 R}{2(a^2 - R^2)}$$